

RECOVERING SPACE CURVES BY MÖBIUS INVARIANTS*

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ABSTRACT: *One subgroup of the Möbius group in the Euclidean 3-space \mathbb{R}^3 is induced by the group of rotations on the unit 3-sphere in the Euclidean 4-space \mathbb{R}^4 via a stereographic projection. We find two functions as Möbius invariants and they determine a Frenet space curve up to a transformation of considered group. We use these invariants to construct space curves and for that purpose we apply an algorithm in the computer system Mathematica.*

KEYWORDS: *Möbius group, Frenet space curves, Stereographic projection*

2010 Math. Subject Classification: *51B10, 53A55*

1 Introduction

The three dimensional unit sphere \mathbb{S}^3 , centered at the origin O and embedded in the four-dimensional Euclidean space \mathbb{R}^4 , is mapped via the stereographic projection π with a center at a point $P \in \mathbb{S}^3$ onto an equatorial hyperplane $\xi \subset \mathbb{R}^4$ through the origin O with a normal vector \overrightarrow{OP} . We consider the hyperplane ξ as a three-dimensional Euclidean space \mathbb{R}^3 . The group of rigid motion on the sphere \mathbb{S}^3 coincides with the group of rotations $SO(4)$ in \mathbb{R}^4 about the origin O . This group induces in \mathbb{R}^3 via the stereographic projection π a subgroup M_0 of the Möbius group $\text{Möb}(3)$ in \mathbb{R}^3 . For that group it is naturally to apply the algebra of quaternions \mathbb{H} , where $\mathbb{H}_0 \subset \mathbb{H}$ is the set of pure quaternions (with zero real parts). If $w \in \mathbb{H}$ is a quaternion then the norm of w is denoted by

*This paper is (partially) supported by the National Scientific Program "Information and Communication Technologies for a Single Digital Market in Science, Education and Security (ICTinSES)", financed by the Ministry of Education and Science.

$\|w\|$, the quaternion $\bar{w} \in \mathbb{H}$ is the conjugate quaternion of w and $\|w\|^2 = w \cdot \bar{w} = \bar{w} \cdot w$. The transformations from the group M_0 can be represented by the quaternion formalism

$$F(q) = (A \cdot q + B) \cdot (B \cdot q + A)^{-1},$$

where $A, B \in \mathbb{H}$, $q \in \mathbb{H}_0$, $\|A\|^2 + \|B\|^2 = 1$ and $A \cdot \bar{B} + B \cdot \bar{A} = 0$, identifying \mathbb{H}_0 with \mathbb{R}^3 (Theorem 3 in [2]).

It is well-known that the Euclidean curvature \varkappa and the Euclidean torsion τ determine the space curves up to orientation preserving motions in \mathbb{R}^3 . This fundamental theorem in the classical curve theory is extended to the orientation preserving similarity group in [1]. Recovering a space curve by a pair of real functions up to an orientation preserving motion or a similarity transformation is based on the mentioned theorems. The Euclidean geometry of curves on $S^3 \subset \mathbb{R}^4$ is carried naturally on \mathbb{R}^3 by the stereographic projection π and it can be interpreted with respect to the subgroup $M_0 \subset \text{Möb}(3)$ in \mathbb{R}^3 . The differential geometry of curves with respect to the Möbius group is developed by Udo Hertrich-Jeromin in [7]. In the presented paper it is used a different approach. Relations between the Euclidean differential-geometric invariants of curves on the sphere S^3 and the corresponding curves in \mathbb{R}^3 via the stereographic projection π are found. These relations determine invariants of curves with respect to the group M_0 . They have an essential roll of recovering the curves in \mathbb{R}^3 up to transformations from the considered group. The case of plane curves, recovering by one Möbuis invariant, is explored in [3] and [4].

2 Preliminaries

Let $K = O\vec{e}_1\vec{e}_2\vec{e}_3\vec{e}_4$ be a right-handed Cartesian coordinate system in \mathbb{R}^4 . For any two vectors $\mathbf{x} = (x^1, x^2, x^3, x^4)$ and $\mathbf{y} = (y^1, y^2, y^3, y^4)$ in \mathbb{R}^4 the canonical inner product $\langle \mathbf{x}, \mathbf{y} \rangle$ is defined by $\langle \mathbf{x}, \mathbf{y} \rangle = \sum_{i=1}^4 x^i y^i$. We denote the norm of the vector \mathbf{x} by $\|\mathbf{x}\| = \sqrt{\langle \mathbf{x}, \mathbf{x} \rangle}$. The inner product and the norm of vectors in \mathbb{R}^3 are determined in a

similar way.

Let $\mathbb{S}^3 = \{(x^1, x^2, x^3, x^4) \in \mathbb{R}^4 \mid (x^1)^2 + (x^2)^2 + (x^3)^2 + (x^4)^2 = 1\}$ be the unit sphere in \mathbb{R}^4 , centered at the origin O , and let $l = (x^1, x^2, x^3, x^4)$, $x^1 \neq 1$ be the position vector of an arbitrary point on \mathbb{S}^3 , different from the pole $P(1, 0, 0, 0)$. Denoting by (u^1, u^2, u^3) the coordinates of any position vector u of an arbitrary point in the equatorial hyperplane $\xi \equiv \mathbb{R}^3$ with respect to the Cartesian coordinate system $K' = O\vec{e}'_1\vec{e}'_2\vec{e}'_3$ in \mathbb{R}^3 , where $\vec{e}'_1 = \vec{e}_2$, $\vec{e}'_2 = \vec{e}_3$, $\vec{e}'_3 = \vec{e}_4$, the stereographic map $\pi : \mathbb{S}^3 \setminus \{P\} \rightarrow \mathbb{R}^3$ is given by the equalities $u^i = \frac{x^{i+1}}{1-x^1}$, $i = 1, 2, 3$.

For the reverse map π^{-1} we have the equalities

$$\begin{aligned} x^1 &= \frac{(u^1)^2 + (u^2)^2 + (u^3)^2 - 1}{(u^1)^2 + (u^2)^2 + (u^3)^2 + 1}, \\ x^i &= \frac{2u^{i-1}}{(u^1)^2 + (u^2)^2 + (u^3)^2 + 1}, \quad i = 2, 3, 4. \end{aligned}$$

Making the infinity ∞ correspond to the pole P , the stereographic projection π and the reverse map π^{-1} are bijective conformal maps.

We consider the parametric equations of the sphere \mathbb{S}^3 in the form

$$l(u^1, u^2, u^3) = \left\{ \frac{\|u\|^2 - 1}{\|u\|^2 + 1}, \frac{2u^1}{\|u\|^2 + 1}, \frac{2u^2}{\|u\|^2 + 1}, \frac{2u^3}{\|u\|^2 + 1} \right\}.$$

The standard orthogonal tangential frame field $\left\{ \frac{\partial l}{\partial u^1}, \frac{\partial l}{\partial u^2}, \frac{\partial l}{\partial u^3} \right\}$ at any point of the sphere \mathbb{S}^3 , except the point P , is given by

$$\begin{aligned} l_1 = \frac{\partial l}{\partial u^1} &= \mu \cdot \left\{ u^1, \frac{1 - (u^1)^2 + (u^2)^2 + (u^3)^2}{2}, -u^1 u^2, -u^1 u^3 \right\} \\ l_2 = \frac{\partial l}{\partial u^2} &= \mu \cdot \left\{ u^2, -u^1 u^2, \frac{1 + (u^1)^2 - (u^2)^2 + (u^3)^2}{2}, -u^2 u^3 \right\} \\ l_3 = \frac{\partial l}{\partial u^3} &= \mu \cdot \left\{ u^3, -u^1 u^3, -u^2 u^3, \frac{1 + (u^1)^2 + (u^2)^2 - (u^3)^2}{2} \right\}, \end{aligned}$$

where $\mu = \frac{4}{(1 + \|\mathbf{u}\|^2)^2}$ and $l_1^2 = l_2^2 = l_3^2 = \mu$.

A regular space curve in \mathbb{R}^3 with a nowhere vanishing Euclidean curvature is called a Frenet space curve. For a given Frenet space curve $c : \mathbf{u}(s) = (u^1(s), u^2(s), u^3(s))$ in \mathbb{R}^3 , parameterized by an arc-length parameter s , we denote by $'$ the differentiation with respect to s . Applying the Frenet equations for the derivatives (see [5]), it is easy to prove the equalities

$$(1) \quad \begin{aligned} \langle \mathbf{u}', \mathbf{u}''' \rangle &= - \langle \mathbf{u}'', \mathbf{u}'' \rangle = -\varkappa^2, \\ \langle \mathbf{u}'', \mathbf{u}''' \rangle &= \left(\frac{\mathbf{u}''^2}{2} \right)' = \varkappa \varkappa', \quad \langle \mathbf{u}' \times \mathbf{u}''', \mathbf{u}'' \rangle = -\varkappa^2 \tau, \end{aligned}$$

where \varkappa and τ are the Euclidean curvature and the Euclidean torsion of c , respectively.

Lemma 2.1. *Let $\mathbf{v} = \mathbf{v}(s)$ be a vector function in \mathbb{R}^3 with coordinate functions $v^k(s) = \sum_{i,j=1}^3 \Gamma_{ij}^k u^i u^j$, $k = 1, 2, 3$, where Γ_{ij}^k are the Christoffel symbols for the sphere \mathbb{S}^3 and $u^i = u^i(s)$, $i = 1, 2, 3$ are the coordinate functions of the unit speed curve $c : \mathbf{u} = \mathbf{u}(s)$. Then*

$$\mathbf{v} = \sqrt{\mu} \cdot \mathbf{u} - \mu \cdot \left(\frac{1}{\mu} \right)' \cdot \mathbf{u}'$$

and

$$\begin{aligned} \langle \mathbf{v}, \mathbf{u}''' \rangle &= \varkappa^2 \mu \left(\frac{1}{\mu} \right)' - \sqrt{\mu} \left(\frac{\sqrt{\mu} + \frac{\mu^2}{4} \left(\frac{1}{\mu} \right)^2 - \frac{\mu}{2} \left(\frac{1}{\mu} \right)''}{\sqrt{\mu}} \right)', \\ \mathbf{v}^2 &= \mu \mathbf{u}^2, \quad \langle \mathbf{v}, \mathbf{u}' \rangle = -\frac{\mu}{2} \left(\frac{1}{\mu} \right)', \quad \langle \mathbf{u}' \times \mathbf{u}''', \mathbf{v} \rangle = \sqrt{\mu} \langle \mathbf{u} \times \mathbf{u}', \mathbf{u}'' \rangle', \end{aligned}$$

$$\begin{aligned} \langle \mathbf{v}, \mathbf{u}'' \rangle &= -\sqrt{\mu} - \frac{\mu^2}{4} \left(\frac{1}{\mu} \right)^2 + \frac{\mu}{2} \left(\frac{1}{\mu} \right)'', \\ \langle \mathbf{u}' \times \mathbf{u}'', \mathbf{v} \rangle &= \mu \langle \mathbf{u} \times \mathbf{u}', \mathbf{u}'' \rangle. \end{aligned}$$

Proof. The Christoffel symbols Γ_{ij}^k for the sphere \mathbb{S}^3 are given by

$$(2) \quad \Gamma_{ij}^k = \Gamma_{ji}^k = \begin{cases} -\sqrt{\mu} \cdot u^j, & k = i; \\ \sqrt{\mu} \cdot u^k, & k \neq i = j; \\ 0, & k \neq i \neq j. \end{cases}$$

Hence,

$$\begin{aligned} \sum_{i,j=1}^3 \Gamma_{ij}^1 u^i u^j &= \Gamma_{11}^1 (u^1)^2 + 2\Gamma_{12}^1 u^1 u^2 + 2\Gamma_{13}^1 u^1 u^3 + \Gamma_{22}^1 (u^2)^2 + \\ &+ \Gamma_{33}^1 (u^3)^2 = -\sqrt{\mu} u^1 (u^1)^2 - 2\sqrt{\mu} u^2 u^1 u^2 - 2\sqrt{\mu} u^3 u^1 u^3 + \\ &+ \sqrt{\mu} u^1 (u^2)^2 + \sqrt{\mu} u^1 (u^3)^2 = -2\sqrt{\mu} u^1 (u^1 u^1 + u^2 u^2 + u^3 u^3) + \\ &+ \underbrace{\sqrt{\mu} u^1 ((u^1)^2 + (u^2)^2 + (u^3)^2)}_{=1} = -\sqrt{\mu} (u^2)' u^1 + \sqrt{\mu} u^1. \end{aligned}$$

Similarly, $\sum_{i,j=1}^3 \Gamma_{ij}^2 u^i u^j = -\sqrt{\mu} (u^2)' u^2 + \sqrt{\mu} u^2$ and $\sum_{i,j=1}^3 \Gamma_{ij}^3 u^i u^j = -\sqrt{\mu} (u^2)' u^3 + \sqrt{\mu} u^3$. Since $(u^2)' = \|\mathbf{u}\|^{2'} = \left(\frac{2}{\sqrt{\mu}} - 1 \right)' = \sqrt{\mu} \left(\frac{1}{\mu} \right)'$, we get the expression of the vector function \mathbf{v} . The next scalar products in the statement of the lemma are obtained by (1) and

$$\langle \mathbf{u}, \mathbf{u}' \rangle = \frac{\sqrt{\mu}}{2} \left(\frac{1}{\mu} \right)' \text{ after the necessary differentiations.} \quad \square$$

The proofs of the next lemmas are routine, replacing the Christoffel symbols by (2).

Lemma 2.2. Let $\mathbf{w} = \mathbf{w}(s)$ be a vector function in \mathbb{R}^3 with coordinate functions $w^k(s) = \sum_{i,j=1}^3 \Gamma_{ij}^k u^i u^{j'}$, $k = 1, 2, 3$, where Γ_{ij}^k are the Christoffel symbols for the sphere \mathbb{S}^3 and $u^i = u^i(s)$, $i = 1, 2, 3$ are the coordinate functions of the unit speed curve $c : \mathbf{u} = \mathbf{u}(s)$. Then

$$\mathbf{w} = -\frac{\mu}{2} \left(\frac{1}{\mu} \right)' \cdot \mathbf{u}'' + \frac{1}{4} \left(4\sqrt{\mu} + \mu^2 \left(\frac{1}{\mu} \right)'^2 - 2\mu \left(\frac{1}{\mu} \right)'' \right) \cdot \mathbf{u}'$$

and

$$\begin{aligned} \langle \mathbf{w}, \mathbf{u}'' \rangle &= -\frac{\mu}{2} \left(\frac{1}{\mu} \right)' \mu^2, \quad \langle \mathbf{w}, \mathbf{v} \rangle = 0, \\ \langle \mathbf{w}, \mathbf{u}' \rangle &= -\langle \mathbf{v}, \mathbf{u}'' \rangle = \sqrt{\mu} + \frac{\mu^2}{4} \left(\frac{1}{\mu} \right)'^2 - \frac{\mu}{2} \left(\frac{1}{\mu} \right)'' . \end{aligned}$$

Lemma 2.3. Let $\boldsymbol{\eta} = \boldsymbol{\eta}(s)$ be a vector function in \mathbb{R}^3 with coordinate functions $\eta^k(s) = \sum_{i,j=1}^3 \Gamma_{ij}^k u^i (u' \times u'')^j$, $k = 1, 2, 3$, where Γ_{ij}^k are the Christoffel symbols for the sphere \mathbb{S}^3 , $u^i = u^i(s)$, $i = 1, 2, 3$ are the coordinate functions of the unit speed curve $c : \mathbf{u} = \mathbf{u}(s)$ and $(u' \times u'')^j$, $j = 1, 2, 3$ are the coordinate functions of the vector product $\mathbf{u}' \times \mathbf{u}''$ in \mathbb{R}^3 . Then

$$\boldsymbol{\eta} = -\frac{\mu}{2} \left(\frac{1}{\mu} \right)' \cdot (\mathbf{u}' \times \mathbf{u}'') - \sqrt{\mu} \langle \mathbf{u} \times \mathbf{u}', \mathbf{u}'' \rangle \cdot \mathbf{u}' .$$

and

$$\begin{aligned} \langle \boldsymbol{\eta}, \mathbf{u}'' \rangle &= \langle \boldsymbol{\eta}, \mathbf{v} \rangle = 0, \\ \langle \boldsymbol{\eta}, \mathbf{u}' \rangle &= -\langle \mathbf{u}' \times \mathbf{u}'', \mathbf{v} \rangle = -\sqrt{\mu} \langle \mathbf{u} \times \mathbf{u}', \mathbf{u}'' \rangle . \end{aligned}$$

Let c be a Frenet space curve in \mathbb{R}^3 . We denote by γ the stereographic pre-image of c on the sphere \mathbb{S}^3 , i.e. $\pi(\gamma) = c$. Let $\mathbf{l} = \mathbf{l}(\sigma)$, $\sigma \in J$ be the vector parametric equation of γ , parameterized by an arc-length

parameter σ . We have that $l^2 = 1$, $l'^2 = 1$ and $\langle l, l' \rangle = 0$, $\langle l, l'' \rangle = -1$, $\langle l, l''' \rangle = 0$. The vector field $l + l''$ is orthogonal to l and $(l + l'')^2 = l'^2 - 1 \geq 0$. A curve γ is said to be regular if $(l + l'')^2 = l'^2 - 1 > 0$ for every $\sigma \in J$. A moving frame $ltnb$ at any point of a regular curve γ on \mathbb{S}^3 is introduced as follows: $t = l'$, $n = \frac{l + l''}{\|l + l''\|}$ and b is the unique vector, determined by the condition $ltnb$ to be a right-handed orthonormal frame. The function $\tilde{\kappa} = \|l + l''\| = \sqrt{l'^2 - 1} > 0$ is called a spherical curvature of γ . The vector field

$$n' + \tilde{\kappa}t = -\frac{\tilde{\kappa}'}{\tilde{\kappa}^2}(l + l'') + \frac{1 + \tilde{\kappa}^2}{\tilde{\kappa}}l' + \frac{1}{\tilde{\kappa}}l'''$$

has the direction of the vector b . So, we can put $n' + \tilde{\kappa}t = \tilde{\tau}b$. The function $\tilde{\tau}$ is called a spherical torsion of γ . For the moving frame $ltnb$ we obtain the following Frenet type formulas

$$l' = t, \quad t' = -l + \tilde{\kappa}n, \quad n' = -\tilde{\kappa}t + \tilde{\tau}b, \quad b' = -\tilde{\tau}n.$$

If we denote by ∇_t the covariant differentiation along the curve γ on \mathbb{S}^3 then $\nabla_t t = \tilde{\kappa}n$, $\nabla_t n = -\tilde{\kappa}t + \tilde{\tau}b$ and $\nabla_t b = -\tilde{\tau}n$. Hence, $\tilde{\kappa} = \|\nabla_t t\|$ and $\tilde{\tau} = -\langle \nabla_t b, n \rangle$ (see [6]).

3 Main theorem

The fundamental theorem in the space curve theory states that the Euclidean curvature κ and the Euclidean torsion τ determine the space curve, parameterized by an arc-length parameter, up to a rigid motion in \mathbb{R}^3 . Also, the spherical curvature $\tilde{\kappa}$ and the spherical torsion $\tilde{\tau}$ define the curve, lying on the sphere $\mathbb{S}^3 \subset \mathbb{R}^4$, up to a rotation, preserving the sphere \mathbb{S}^3 . Relations between the invariants $\kappa, \tau, \tilde{\kappa}, \tilde{\tau}$ of the corresponding curves via the stereographic projection π will define any Frenet space curve up to transformation of the group M_0 .

Theorem 3.1. *Let c be a Frenet space curve in \mathbb{R}^3 , parameterized by an arc-length parameter s , and let γ be its stereographic pre-image on*

the sphere \mathbb{S}^3 , i.e. $\pi(\gamma) = c$. If $\varkappa = \varkappa(s)$ and $\tau = \tau(s)$ are the Euclidean curvature and the Euclidean torsion of the curve c and $\tilde{\varkappa} = \tilde{\varkappa}(s)$ and $\tilde{\tau} = \tilde{\tau}(s)$ are the spherical curvature and the spherical torsion of the curve γ then

$$(3) \quad \tilde{\varkappa}^2 = \frac{1}{\mu} \varkappa^2 - 1 - \frac{3\mu}{4} \left(\frac{1}{\mu}\right)^{\prime 2} + \left(\frac{1}{\mu}\right)''$$

$$(4) \quad \tilde{\tau} = \frac{\cos \theta}{\mu \tilde{\varkappa}} \varkappa \tau - \frac{\sin \theta}{\mu \tilde{\varkappa}} \varkappa',$$

where

$$(5) \quad \cos \theta = \frac{1}{\sqrt{\mu \varkappa \tilde{\varkappa}}} \left(\varkappa^2 - \frac{1}{4} \mu^2 \left(\frac{1}{\mu}\right)^{\prime 2} + \frac{1}{2} \mu \left(\frac{1}{\mu}\right)'' - \sqrt{\mu} \right).$$

Proof. Let $\mathbf{l} = \mathbf{l}(s)$ be a vector parametric equation of γ and $\sigma = \sigma(s)$ be its arc-length function. If $c : \mathbf{u}(s) = (u^1(s), u^2(s), u^3(s))$ then $\dot{\mathbf{l}} = \frac{d\mathbf{l}}{ds} = \sum_{i=1}^3 (u^i)' \mathbf{l}_i$. The form $d\sigma = \sqrt{\mu} ds$ is the linear element of the curve γ on the sphere \mathbb{S}^3 . Thus, the unit tangent vector field is represented by $\mathbf{t} = \frac{1}{\sqrt{\mu}} \dot{\mathbf{l}}$. Applying the covariant differentiation $\nabla_{\mathbf{t}}$ we obtain that

$$(6) \quad \nabla_{\mathbf{t}} \mathbf{t} = \sum_{k=1}^3 \left(\frac{1}{\mu} (u^k)'' + \frac{1}{\mu} v^k + \frac{1}{2} \left(\frac{1}{\mu}\right)' (u^k)' \right) \mathbf{l}_k,$$

where $v^k = v^k(s)$, $k = 1, 2, 3$ are the coordinate functions of the vector function \mathbf{v} , defined in Lemma 2.1. Hence, $\tilde{\varkappa}^2 =$

$$= \|\nabla_{\mathbf{t}} \mathbf{t}\|^2 = \frac{1}{\mu} \varkappa^2 + \frac{1}{\mu} v^2 + \frac{\mu}{4} \left(\frac{1}{\mu}\right)^{\prime 2} + \frac{2}{\mu} \langle \mathbf{u}'', \mathbf{v} \rangle + \left(\frac{1}{\mu}\right)' \langle \mathbf{v}, \mathbf{u}' \rangle$$

and replacing with the equalities in Lemma 2.1 we get (3). The stereographic projection π and the reverse map π^{-1} are conformal maps and the differential π_*^{-1} of the reverse map π^{-1} is a similarity transformation between the corresponding tangent spaces with a coefficient $\sqrt{\mu}$. If $T = \mathbf{u}'$, $N = \frac{\mathbf{u}''}{\varkappa}$, $B = T \times N = \frac{\mathbf{u}' \times \mathbf{u}''}{\varkappa}$ are the Frenet frame vector field along the curve c then $\pi_*^{-1}(N) = \frac{1}{\varkappa} \sum_{i=1}^3 (u^i)'' l_i$ and $\pi_*^{-1}(B) = \frac{1}{\varkappa} \sum_{i=1}^3 (u' \times u'')^i l_i$. Vectors $\pi_*^{-1}(N)$ and $\pi_*^{-1}(B)$, lying in the plane, determined by vectors \mathbf{n} and \mathbf{b} , are orthogonal and let $\theta = \angle(\mathbf{n}, \pi_*^{-1}(N))$. Then $\mathbf{n} = \frac{\cos \theta}{\sqrt{\mu}} \pi_*^{-1}(N) + \frac{\sin \theta}{\sqrt{\mu}} \pi_*^{-1}(B)$. Hence,

$$\cos \theta = \frac{1}{\sqrt{\mu}} \langle \mathbf{n}, \pi_*^{-1}(N) \rangle = \frac{1}{\tilde{\varkappa} \varkappa \sqrt{\mu}} \langle \nabla_t \mathbf{t}, \sum_{i=1}^3 (u^i)'' l_i \rangle$$

and

$$\sin \theta = \frac{1}{\sqrt{\mu}} \langle \mathbf{n}, \pi_*^{-1}(B) \rangle = \frac{1}{\tilde{\varkappa} \varkappa \sqrt{\mu}} \langle \nabla_t \mathbf{t}, \sum_{i=1}^3 (u' \times u'')^i l_i \rangle.$$

Using (6) we obtain that $\cos \theta = \frac{1}{\tilde{\varkappa} \varkappa \sqrt{\mu}} (\mathbf{u}''^2 + \langle \mathbf{v}, \mathbf{u}'' \rangle) = \frac{1}{\tilde{\varkappa} \varkappa \sqrt{\mu}} (\varkappa^2 + \langle \mathbf{v}, \mathbf{u}'' \rangle)$, whence applying Lemma 2.1 we get (5). Also,

$$(7) \quad \sin \theta = \frac{1}{\tilde{\varkappa} \varkappa \sqrt{\mu}} \langle \mathbf{v}, \mathbf{u}' \times \mathbf{u}'' \rangle = \frac{1}{\tilde{\varkappa} \varkappa} \langle \mathbf{u} \times \mathbf{u}', \mathbf{u}'' \rangle.$$

Observing that π is an orientation preserving map in odd dimension, we have that $\mathbf{b} = \frac{-\sin \theta}{\sqrt{\mu}} \pi_*^{-1}(N) + \frac{\cos \theta}{\sqrt{\mu}} \pi_*^{-1}(B) =$

$= -\frac{1}{\varkappa\sqrt{\mu}} \sum_{i=1}^3 (\sin \theta (u^i)'' - \cos \theta (u' \times u'')^i) l_i$. Therefore, using the covariant differentiation technique, we get

$$\begin{aligned} \nabla_t \mathbf{b} &= \frac{1}{\sqrt{\mu}} \nabla_i \mathbf{b} = -\frac{1}{\sqrt{\mu}} \sum_{k=1}^3 \left[\frac{1}{\varkappa\sqrt{\mu}} (\sin \theta (w^k + (u^k)''') - \right. \\ &\left. - \cos \theta (\eta^k + (u' \times u'')^k)) + \left(\frac{\sin \theta}{\varkappa\sqrt{\mu}} \right)' (u^k)'' - \left(\frac{\cos \theta}{\varkappa\sqrt{\mu}} \right)' (u' \times u'')^k \right] l_k. \end{aligned}$$

Thus,

$$\begin{aligned} \tilde{\tau} &= -\langle \nabla_t \mathbf{b}, n \rangle = -\frac{1}{\tilde{\varkappa}} \langle \nabla_t \mathbf{b}, \nabla_t \mathbf{t} \rangle = \frac{\sqrt{\mu}}{\tilde{\varkappa}} \left[\frac{\sin \theta}{\varkappa\mu\sqrt{\mu}} \cdot \right. \\ &\left(\underbrace{\langle w, u'' \rangle + \langle u''', u'' \rangle + \langle w, v \rangle + \langle u''', v \rangle}_A \right) + \frac{\sin \theta}{2\varkappa\sqrt{\mu}} \left(\frac{1}{\mu} \right)' \cdot \\ &\left(\underbrace{\langle w, u' \rangle + \langle u''', u' \rangle}_B \right) - \frac{\cos \theta}{\varkappa\mu\sqrt{\mu}} \cdot \\ &\left(\underbrace{\langle \eta, u'' \rangle + \langle u' \times u''', u'' \rangle + \langle \eta, v \rangle + \langle u' \times u''', v \rangle}_C \right) - \\ &\quad - \frac{\cos \theta}{2\varkappa\sqrt{\mu}} \left(\frac{1}{\mu} \right)' \underbrace{\langle \eta, u' \rangle}_D + \\ &\left. + \frac{1}{\mu} \left(\frac{\sin \theta}{\varkappa\sqrt{\mu}} \right)' \left(\underbrace{\varkappa^2 + \langle u''', v \rangle}_E \right) - \frac{1}{\mu} \left(\frac{\cos \theta}{\varkappa\sqrt{\mu}} \right)' \underbrace{\langle u' \times u''', v \rangle}_F \right] = \end{aligned}$$

$$= \frac{1}{\tilde{\varkappa}} \left[\frac{\sin \theta}{\varkappa \mu} .A + \frac{\sin \theta}{2\varkappa} \left(\frac{1}{\mu} \right)' .B - \frac{\cos \theta}{\varkappa \mu} .C - \frac{\cos \theta}{2\varkappa} \left(\frac{1}{\mu} \right)' .D + \right. \\ \left. + \frac{1}{\sqrt{\mu}} \left(\frac{\sin \theta}{\varkappa \sqrt{\mu}} \right)' .E - \frac{1}{\sqrt{\mu}} \left(\frac{\cos \theta}{\varkappa \sqrt{\mu}} \right)' .F. \right]$$

Applying the lemmas 2.1, 2.2, 2.3 and taking account of the equalities (1), (5), (7) we have that $A = -\varkappa \varkappa' + \sqrt{\mu}(\varkappa \tilde{\varkappa} \cos \theta)'$, $-B = E = \sqrt{\mu} \varkappa \tilde{\varkappa} \cos \theta$, $C = -\varkappa^2 \tau + \sqrt{\mu}(\varkappa \tilde{\varkappa} \sin \theta)'$, $-D = F = \sqrt{\mu} \varkappa \tilde{\varkappa} \sin \theta$. Replacing the found expressions in the equality above we obtain (4) after simple rearranges and simplifications. \square

4 Möbius invariants of space curves

It is obvious that the arc-length function $\sigma = \sigma(s)$ is an invariant with respect to the group M_0 . Let us denote by $S_\sigma(s)$ the Schwarzian derivative of the function $\sigma = \sigma(s)$. Then we have that

$$S_\sigma(s) = \left(\frac{\sigma''}{\sigma'} \right)' - \frac{1}{2} \left(\frac{\sigma''}{\sigma'} \right)^2 = \frac{1}{2} \left[\frac{3}{4} \mu^2 \left(\frac{1}{\mu} \right)'^2 - \mu \left(\frac{1}{\mu} \right)'' \right].$$

If $s = s(\sigma)$ is the reverse function of σ then we get $0 = S_{\sigma \circ s}(\sigma) = S_\sigma(s) \left(\frac{ds}{d\sigma} \right)^2 + S_s(\sigma) \Rightarrow S_\sigma(s) = -\mu S_s(\sigma)$ applying the chain rule for the Schwarzian derivatives. The functions $\tilde{\varkappa}$ and $\tilde{\tau}$ from Theorem 3.1 can be expressed in terms of the arc-length parameter σ . We set $\mathfrak{K}(\sigma) = \tilde{\varkappa}^2(\sigma)$, $\mathfrak{T}(\sigma) = \tilde{\tau}(\sigma)$ and we obtain that

$$(8) \quad \mathfrak{K}(\sigma) = \frac{1}{\mu} \varkappa^2 - 1 + 2S_s(\sigma) \\ \mathfrak{T}(\sigma) = \frac{1}{\mu \sqrt{\mathfrak{K}}} \left(\varkappa \tau \cos \theta - \sqrt{\mu} \frac{d\varkappa}{d\sigma} \sin \theta \right),$$

$$\text{where } \cos \theta = \frac{1}{\varkappa \sqrt{\mu \mathfrak{K}}} \left(\varkappa^2 + \mu S_s(\sigma) + \frac{1}{8\mu} \left(\frac{d\mu}{d\sigma} \right)^2 - \sqrt{\mu} \right).$$

Corollary 4.1. *The functions $\mathfrak{K} = \mathfrak{K}(\sigma)$ and $\mathfrak{T} = \mathfrak{T}(\sigma)$, define by (8), are invariants under the group M_0 .*

Proof. The proof is omitted. □

Theorem 4.2 (Uniqueness theorem). *Let $I \subset \mathbb{R}$ be an open interval and let $c_i : I \rightarrow \mathbb{R}^3$, $i = 1, 2$ be two Frenet space curves, parameterized by the same arc-length parameter σ of their stereographic pre-images γ_i , $i = 1, 2$, respectively, on the sphere \mathbb{S}^3 in \mathbb{R}^4 . Assume that the curves c_1 and c_2 have the same invariants $\mathfrak{K}_i = \mathfrak{K}_i(\sigma)$, $\mathfrak{T}_i = \mathfrak{T}_i(\sigma)$, $i = 1, 2$, defined by (8) for any $\sigma \in I$. Then there exists a transformation $F \in M_0$ such that $c_2 = F(c_1)$.*

Proof. Since $\mathfrak{K}_1 = \mathfrak{K}_2$ and $\mathfrak{T}_1 = \mathfrak{T}_2$ then the spherical stereographic pre-images γ_1 and γ_2 of the curves c_1 and c_2 , respectively, via the stereographic projection π , have the same spherical curvatures and torsions. Therefore, there exists a rotation $\rho \in SO(4)$ such that $\gamma_2 = \rho(\gamma_1)$. Hence, $c_2 = \pi(c_1) = \pi \circ \rho(\gamma_1) = \pi \circ \rho \circ \pi^{-1}(c_1)$, where $F = \pi \circ \rho \circ \pi^{-1} \in M_0$ and the proof is completed. □

Theorem 4.3 (Existence theorem). *Let $f : I \rightarrow \mathbb{R}$, $f > 0$ and $g : I \rightarrow \mathbb{R}$ be given C^∞ -functions, defined on the same interval $I \subset \mathbb{R}$. Let $\mathbf{c}_0 \in \mathbb{R}^3$ and $\mathbf{e}_1^0, \mathbf{e}_2^0, \mathbf{e}_3^0$ be a right-handed orthonormal frame at \mathbf{c}_0 in the Euclidean space \mathbb{R}^3 . There exists a unique Frenet space curve $c : I \rightarrow \mathbb{R}^3$ which satisfies the conditions:*

- (a) *there exists $\sigma_0 \in I$, such that $c(\sigma_0) = \mathbf{c}_0$, and the Frenet frame of c at \mathbf{c}_0 is $\mathbf{e}_1^0, \mathbf{e}_2^0, \mathbf{e}_3^0$;*
- (b) *for any $\sigma \in I$ $\mathfrak{K}(\sigma) = f^2(\sigma)$ and $\mathfrak{T}(\sigma) = g(\sigma)$.*

Proof. Let $\mathbf{l}_0 = \pi^{-1}(\mathbf{c}_0)$ and $\mathbf{t}_0 = \pi_*^{-1}(\mathbf{e}_1^0)$, $\mathbf{n}_0 = \pi_*^{-1}(\mathbf{e}_2^0)$, $\mathbf{b}_0 = \pi_*^{-1}(\mathbf{e}_3^0)$. Let us consider a matrix-valued function $\mathcal{E}(\sigma) = (\mathbf{l}(\sigma), \mathbf{t}(\sigma), \mathbf{n}(\sigma), \mathbf{b}(\sigma))^T$. Solving the system of first order linear differential equations

$$\frac{d}{d\sigma}\mathcal{E} = \mathcal{A}(\sigma)\mathcal{E}, \text{ with a given matrix } \begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & f & 0 \\ 0 & -f & 0 & g \\ 0 & 0 & -g & 0 \end{pmatrix} \text{ and ini-}$$

tial conditions $\mathbf{l}_0, \mathbf{t}_0, \mathbf{n}_0, \mathbf{b}_0$ we obtain a unique solution $\mathcal{E} = \mathcal{E}(\sigma)$, determined for all $\sigma \in I$ and $\mathcal{E}(\sigma_0) = (\mathbf{l}_0, \mathbf{t}_0, \mathbf{n}_0, \mathbf{b}_0)^T$ for some $\sigma_0 \in I$. It is routine to prove that the matrix \mathcal{E} is orthogonal. This means that the vectors $\mathbf{l}(\sigma), \mathbf{t}(\sigma), \mathbf{n}(\sigma), \mathbf{b}(\sigma)$ form an orthogonal frame in \mathbb{R}^4 for any $\sigma \in I$. Let γ be a spherical curve, defined by the vector function $\mathbf{l} = \mathbf{l}(\sigma)$ and let $c = \pi(\gamma)$. It is clear that the conditions (a) and (b) in the statement of the theorem are fulfilled for the space curve c . \square

The proof of the last theorem give us an algorithm of recovering space curves up to a transformation from the group M_0 .

Algorithm. *Recovering space curves by two functions $f = f(\sigma) > 0$ and $g = g(\sigma)$ for any $\sigma \in I$.*

1. Choose initial conditions: $\mathbf{c}_0, \mathbf{e}_1^0, \mathbf{e}_2^0, \mathbf{e}_3^0$;
2. Find the vectors $\mathbf{l}_0 = \pi^{-1}(\mathbf{c}_0), \mathbf{t}_0 = \pi_*^{-1}(\mathbf{e}_1^0), \mathbf{n}_0 = \pi_*^{-1}(\mathbf{e}_2^0), \mathbf{b}_0 = \pi_*^{-1}(\mathbf{e}_3^0)$;

3. Solve the differential equation $\frac{d}{d\sigma}\mathcal{E} = \mathcal{A}(\sigma)\mathcal{E}$,

$$\text{where } \begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & f & 0 \\ 0 & -f & 0 & g \\ 0 & 0 & -g & 0 \end{pmatrix}, \text{ and initial conditions, determined}$$

in Step 1., for $\mathcal{E}(\sigma_0) = (\mathbf{l}_0, \mathbf{t}_0, \mathbf{n}_0, \mathbf{b}_0)^T$;

4. The curves γ , with vector function $\mathbf{l} = \mathbf{l}(\sigma)$, and $c = \pi(\gamma)$ are found.

The next two examples illustrate the considered conformal invariants of a space curve in \mathbb{R}^3 . For the calculations and visualizations we use the computer system Mathematica.

Example 4.4. *Let*

$$c : \mathbf{u}(\sigma) = \left\{ \cos \left(\sqrt{2} \tan \left(\frac{\sigma}{2} \right) \right), \sin \left(\sqrt{2} \tan \left(\frac{\sigma}{2} \right) \right), \sqrt{2} \tan \left(\frac{\sigma}{2} \right) \right\}$$

be a helix in \mathbb{R}^3 , parameterized by an arc-length parameter σ of its spherical stereographic pre-image. Then $\mathfrak{K}(\sigma) = \frac{1}{4 \cos^4 \left(\frac{\sigma}{2} \right)}$ and $\mathfrak{T}(\sigma) = \frac{\cos \sigma}{1 + \cos \sigma}$.

Example 4.5. *Let $\mathfrak{K}(\sigma) = \sigma^2$, $\mathfrak{T}(\sigma) = 0.6\sigma$, $\mathbf{c}_0 = (0, 0, 0)$, $\mathbf{e}_1^0 = (0, 0, 1)$, $\mathbf{e}_2^0 = (0, 1, 0)$. Then, applying the algorithm above, where $f(\sigma) = \sigma$, and $g(\sigma) = 0.6\sigma$, we obtain the Frenet space curve, depicted in Fig.1.*

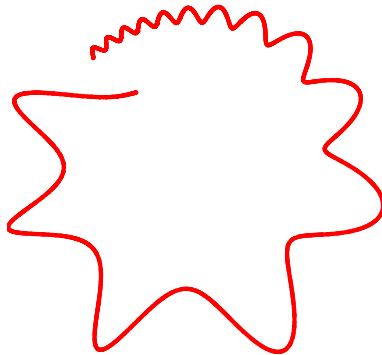


Fig. 1: $\mathfrak{K}(\sigma) = \sigma^2$, $\mathfrak{T}(\sigma) = 0.6\sigma$

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